

1- Basic questions: see lectures

3°) When rod-like molecules lie in a plane, we denote  $\theta$  the angle wrt the direction  $\hat{n}$  and define order param

$$S = \frac{2\langle \cos^2 \theta \rangle - 1}{2 - 1} = 2\langle \cos^2 \theta \rangle - 1$$

4°) Non interacting spins:

$$H = -B \sum_{i=1}^N S_i$$

$$Z = \sum_{S_1, \dots, S_N} e^{-\beta H} = [2 \cosh(\beta B)]^N$$

$$H = \sum_{i=1}^N S_i; \langle n \rangle = \frac{1}{Z} \sum_{S_1, \dots, S_N} M e^{\beta B M}$$

$$\langle n \rangle = \frac{\partial}{\partial (\beta B)} \log Z = - \frac{\partial F}{\partial B} \Big|_T$$

$$\langle n \rangle = N \tanh(\beta B)$$

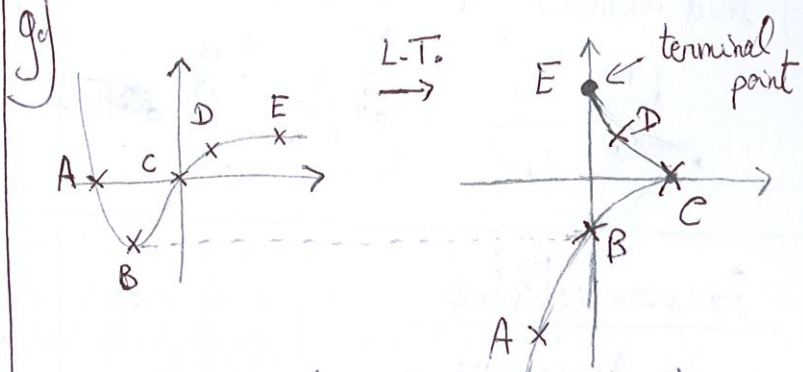
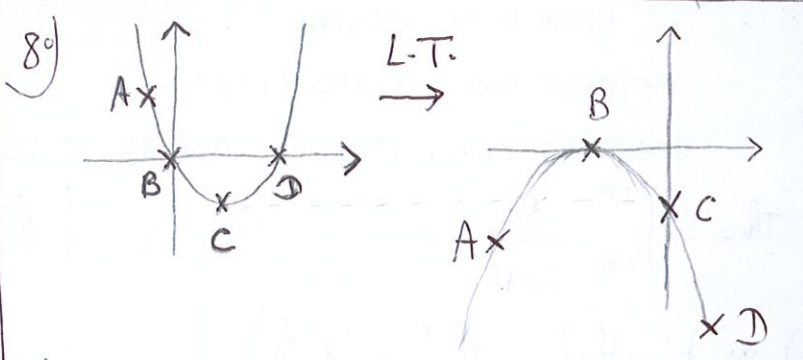
5°)  $\langle n \rangle = - \frac{\partial F}{\partial B} \Big|_T$ , see above;  $F(T, B)$

6°)  $C_B = T \frac{\partial S}{\partial T} \Big|_B$ ;  $dF = -SdT - MdB$

$C_B = -T \frac{\partial^2 F}{\partial T^2} \Big|_B$

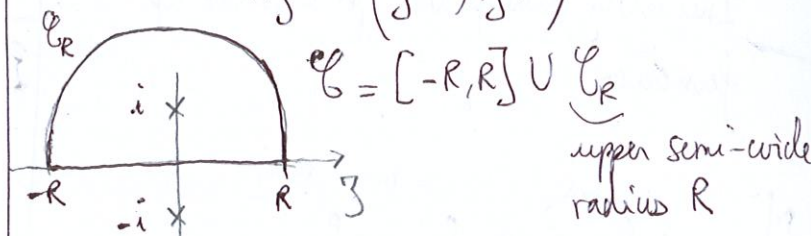
7°)  $C_n = T \frac{\partial S}{\partial T} \Big|_n$ ;  $d\tilde{F} = d(F + \beta n M)$   
 $= -SdT + Bdn$

$C_n = -T \frac{\partial^2 \tilde{F}}{\partial T^2} \Big|_n$ ;  $\tilde{F}(T, n)$   
 $\tilde{F} = F + \beta n M$



this latter graph is not that of a function; due to lack of convexity of the input

10°) a)  $f(z) = \frac{1}{1+z^2} = \frac{1}{(z-i)(z+i)}$  has poles  $\pm i$



Residue theorem:  $\oint_{C_R} f(z) dz = 2i\pi \times \text{residue}$   
 $= 2i\pi \frac{1}{2i} = \pi$

$\int_{-R}^R f + \int_{C_R} f = \pi$   
 can be bound;  $|z^2 + 1| \geq R^2 - 1$   
 $\geq \frac{R^2}{2}$

if R large enough  
 $\Rightarrow \left| \int_{C_R} f(z) dz \right| \leq \left| \int_{C_R} \frac{dz}{|1+z^2|} \right|$   
 $\leq \frac{2}{R^2} \pi R \xrightarrow{R \rightarrow \infty} 0$

Thus:  $\int_{-\infty}^{+\infty} \frac{dx}{1+x^2} = \pi$

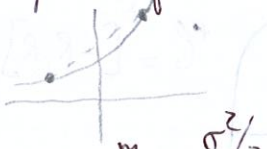
f) Direct method:  $th' = 1 + th^2$   
 $\int_{-\infty}^{+\infty} \frac{dx}{1+x^2} = \left[ \operatorname{arctg} x \right]_{-\infty}^{+\infty} = \pi$

2- Gaussian calculus

$x \rightarrow g(m, \sigma)$

1)  $\langle e^x \rangle = e^m e^{\sigma^2/2}$   
 $e^{\langle x \rangle} = e^m \langle \langle e^x \rangle \rangle$

Inequality always true for a convex up function



2)  $\langle e^{-x} \rangle = e^{-m} e^{\sigma^2/2}$

3)  $p(x) = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-x^2/2\sigma^2}$ ;  $\frac{dp}{dx} = -\frac{x}{\sigma^2} p(x)$

$\langle \frac{df}{dx} \rangle = \int \frac{df}{dx} p(x) dx$   
 $= \left[ f(x)p(x) \right]_{-\infty}^{+\infty} + \frac{1}{\sigma^2} \int x f(x) p(x) dx$

$\Rightarrow \langle x f(x) \rangle = \sigma^2 \langle \frac{df}{dx} \rangle$

4) Take  $f(x) = x^3$ :  $\langle x^4 \rangle = 3\sigma^4 = 3\langle x^2 \rangle^2$   
 $f(x) = x^5$ :  $\langle x^6 \rangle = 5\sigma^2 \langle x^4 \rangle$   
 $= 5 \times 3 \times \sigma^6$

5) Hence  $\langle x^n \rangle = 0$  for  $n$  odd  
 $\langle x^n \rangle = (n-1)(n-3)\dots 1$  for  $n$  even

6)  $e^x = \sum_{n=0}^{\infty} \frac{x^n}{n!}$   
 $\langle e^x \rangle = \sum_{p=0}^{\infty} \frac{\langle x^{2p} \rangle}{(2p)!} =$   
 $= \sum_{p=0}^{\infty} \frac{(2p-1)(2p-3)\dots 1}{(2p)!} \sigma^{2p}$   
 $= \sum_{p=0}^{\infty} \frac{\sigma^{2p}}{(2p)(2p-2)\dots 2}$   
 $= \sum_{p=0}^{\infty} \frac{(\sigma^2/2)^p}{p!} = e^{\sigma^2/2}$ ; OK

3- Magnetic domain walls

BC can be enforced by a strong negative B field for  $x \rightarrow -\infty$  / positive for  $x \rightarrow +\infty$ .

Extremality:  $F = \int L(m, \frac{dm}{dx}) dx$

$\frac{\partial L}{\partial m} = \frac{d}{dx} \left( \frac{\partial L}{\partial \left( \frac{dm}{dx} \right)} \right)$

$\Rightarrow -m + m^3 = \frac{d^2 m}{dx^2}$  Solve... yields  $m^*$

$\Rightarrow \left( \frac{dm}{dx} \right)^2 = \frac{m^4}{2} - m^2 + \frac{1}{2} = \frac{1}{2} (m^2 - 1)^2$

$\frac{dm}{dx} = \frac{1}{\sqrt{2}} (1 - m^2)^{1/2}$

$\Rightarrow m^*(x) = \operatorname{th} \left( \frac{x}{\sqrt{2}} \right)$